

Schemes in Conservation Form

Definition (Conservative Scheme)

A conservation law scheme is called **conservative** iff it can be written as

$$u_{j,l}^{n+1} = u_{j,l}^n - \frac{\Delta t}{\Delta x} \left(f_{j+\frac{1}{2}}^* \left(\vec{u}_l^n \right) - f_{j-\frac{1}{2}}^* \left(\vec{u}_l^n \right) \right)$$

where f^* ...

- Lipschitz-continuous

- "consistency": $f^*(u, u, u, \dots, u) = f(u)$

Theorem (Lax-Wendroff)

If the solution $\{u_{j,l}\}$ to a conservative scheme converges (as $\Delta t, \Delta x \rightarrow 0$) boundedly almost everywhere to a function $u(x, t)$, then u is a weak solution of the conservation law

Lax-Wendroff Theorem: Proof

$\Delta_i^\pm \rightarrow$ applied on index axis i .

Summation by parts: With $\Delta^+ a_k = a_{k+1} - a_k$ and $\Delta^- a_k = a_k - a_{k-1}$:

$$\int_{\omega'} u' - \int_{\omega} u' = \sum_{k=1}^N a_k (\Delta^- \varphi_k) + \sum_{k=1}^N \varphi_k (\Delta^+ a_k) = -a_1 \varphi_0 + \varphi_N a_{N+1}.$$

Let $\varphi_{j,l} = \varphi(x_{j,l}, t_l)$ $\varphi \in C_b^1$ ("compact support").

$$0 = \sum_{l=1}^{\infty} \sum_j \left(\frac{\Delta_x^+ u_{j,l}}{h_x} + \frac{\Delta_t^+ \varphi_{j-\frac{1}{2}}}{h_t} \right) \varphi_{j,l} \Big|_{h_x h_t} \xrightarrow{\text{by cons.}} 0$$

$$= - \sum_{l=1}^{\infty} \sum_j \left(\frac{(\Delta_t^- \varphi_{j,l}) u_{j,l}}{h_t} + \frac{\Delta_x^- \varphi_{j,l}}{h_x} f_{j-\frac{1}{2}}^* \right) h_x h_t - \sum_j u_{j,1} \varphi_{j,0} h_x$$

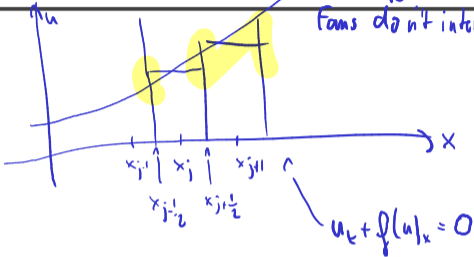
DCT \rightarrow

$$(h_x, h_t \rightarrow 0) \int_0^{\infty} \int_{-\infty}^{\infty} \varphi_t u + \varphi_x f(u) dx dt - \int_{-\infty}^{\infty} u(x,0) \varphi(x,0) dx$$

Finite Volume Schemes

Finite volume: Idea?

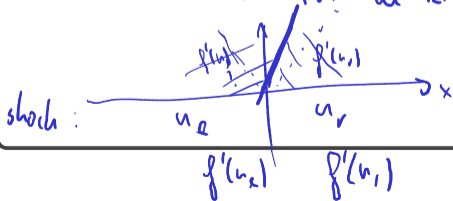
Intuitively; choose Δt so that Riemann fans don't intersect.



$$\bar{u}_j = \frac{1}{\Delta x} \int_{x_{j-1/2}}^{x_{j+1/2}} u(x) dx$$



Idea behind FV: all Riemann problems all the time.



Developing Finite Volume

$$\int_{t_\ell}^{t^{\ell+1}} \int_{x_{j-1/2}}^{x_{j+1/2}} (u_t + f(u)_x) dx dt = 0$$

Handwritten derivation of the finite volume method:

$$\frac{1}{h_x} \int_{x_{j-1/2}}^{x_{j+1/2}} u^{\ell+1}(x) dx - \frac{1}{h_x} \int_{x_{j-1/2}}^{x_{j+1/2}} u_\ell(x) dx + \frac{1}{h_x} \int_{t_\ell}^{t^{\ell+1}} f(u_{j+1/2}(t)) dt - \frac{1}{h_x} \int_{t_\ell}^{t^{\ell+1}} f(u_{j-1/2}(t)) dt = 0$$

Diagram illustrating the control volume (CV) centered at x_j with width h_x . The CV is bounded by $x_{j-1/2}$ and $x_{j+1/2}$. The fluxes at the boundaries are $f(u_{j+1/2})$ and $f(u_{j-1/2})$.

$$\bar{u}_{j,\ell+1} - \bar{u}_{j,\ell} + \frac{1}{h_x} \int_{t_\ell}^{t^{\ell+1}} f(u_{j+1/2}(t)) dt - \frac{1}{h_x} \int_{t_\ell}^{t^{\ell+1}} f(u_{j-1/2}(t)) dt = 0$$

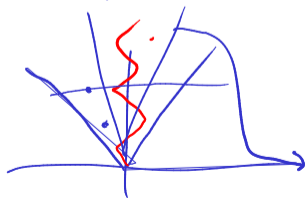
Flux Integrals?

$$\frac{1}{h_x} \int_{t_l}^{t_{l+1}} f(u_{j+1/2}) dt? \quad \leftarrow \text{Halp?}$$

Observes Solutions to Riemann problems are self-similar.

$$\begin{array}{l} T: \quad \bar{x} = ax \quad E = at \\ u_t + f(u)_x = 0 \quad \rightsquigarrow \quad ? u_{\bar{x}} \quad + ? f(u)_{\bar{x}} = 0 \end{array}$$

↳ Change of variables T leaves f c. law invariant.



just straight lines

IC
⇒ Riemann problem

$$u(x,t) = U(x/t)$$

The Godunov Scheme

Altogether:

$$\bar{u}_{j,\ell+1} = \bar{u}_{j,\ell} - \frac{h_t}{h_x} (f(u_{j+1/2,\ell}) - f(u_{j-1/2,\ell})).$$

Overall algorithm?

Heuristic time step restriction?