

Today 2/1

$$Z(x,t) = z_0 e^{i(kx - \omega t)}$$

$= z_0 e^{i2\pi\left(\frac{x}{\lambda} - \frac{t}{T}\right)}$

$z_0 =$ amplitude

$k =$ wave #

$\omega =$ ang frequency

$T =$ period

$= \frac{1}{\nu} \leftarrow$ frequency

$$\omega = 2\pi\nu$$

$\lambda =$ wave length

$$= \frac{2\pi}{k}$$

Main Tool

$$\text{let } z = z_0 e^{i(kx - \omega t)}$$

$$Lz = \lambda z$$

for any linear operator L .

$$L = \partial_t : \partial_t z = -i\omega z$$

also

$$L = \partial_x : \partial_x z = ikz$$

$$z \text{ solve } Lz = 0 \text{ iff } \lambda = 0$$

$\rightarrow \lambda = 0$ is called the dispersion relation.

Example

$$L = \partial_t + a \partial_x$$

$$\rightarrow Lz = -i\omega z + aikz$$

$$= (-i\omega + aik)z$$

$$= \underbrace{i(ak - \omega)}_{} z$$

dispersion relation: $\searrow (k, \omega)$

$$\lambda \equiv 0$$

$$\Rightarrow i(ak - \omega) = 0$$

$$\Rightarrow \omega = ak$$

Example

$$L = \cancel{u_t + a u_x - D u_{xx}}$$
$$= \partial_t + a \partial_x - D \partial_{xx}$$

$$\rightarrow Lz = -i\omega z + aikz - D(ik)^2 z$$
$$= (-i\omega + aik + Dk^2) z$$

$$\rightarrow \boxed{\omega = ak - iDk^2}$$

take $z = z_0 e^{i(kx - \omega t)}$

let $\omega(k) = \underbrace{\alpha(k)}_{\text{real}} + i \underbrace{\beta(k)}_{\text{imag}}$

$= z_0 e^{i(kx - \alpha t - i\beta t)}$

$= z e^{\beta t} e^{i(kx - \alpha t)}$

↑
dissipative

if $\alpha(k) = ak$
then $e^{i(kx - \alpha t)}$ is a fixed
wave at a
fixed velocity

$$\Rightarrow v_{ph} = \frac{\text{Re}(\omega(k))}{k} = \frac{\alpha(k)}{k}$$

if v_{ph} is constant,
then no dispersion
if v_{ph} is non-constant,
then dispersive.

$$\rightarrow z_{j,l} = z_0 e^{i(k_x x - \omega t)}$$

\uparrow \uparrow
 $j k_x$ $l \omega t$

$$= z_0 e^{i(K_j k_x - \omega l \omega t)}$$

Example ETBS

$$u_{j,l,t+1} = \gamma u_{j-1,l} + (1-\gamma) u_{j,l}$$

let $u = z$:

~~$$z_{j,l} e^{-i\omega l \omega t} = \gamma z_{j-1,l} e^{-i k k_x} + (1-\gamma) z_{j,l}$$~~

$$\Rightarrow e^{-i\omega l \omega t} = \underbrace{\gamma e^{-i k k_x} + (1-\gamma)}_{\text{amp factor}}$$

$$e^{-i\omega t} = \underbrace{\gamma e^{-ikhx} + 1 - \gamma}$$

$S(kh_x) =$ symbol
or
amp factor

stability: if $|S(kh_x)| \leq 1$
then stable.

$$kh_x = ? = \frac{2\pi h_x}{\lambda}$$

$$\omega = \omega(kh_x)$$

$$\begin{aligned} \text{look at } s &= \text{complex} \\ &= |s| e^{i\phi} \quad (\text{some } \phi) \\ &= e^{\ln|s| + i\phi} \end{aligned}$$

$$\text{want } e^{-i\omega h t} = e^{\ln|s| + i\phi}$$

goal: find ω

$$-i\omega h t = \ln|s| + i\phi$$

$$\omega h t = i \ln|s| - \phi$$

$$\omega =$$

$$\frac{i \ln|s| - \phi}{h t}$$

back to $z_{j\ell}$

$$z_{j\ell} = z_0 e^{i(k_j h_x - \omega \ell h_t)}$$

↑ put it here

$$= z_0 e^{i(k_j h_x - \frac{i|\ln|s| - \phi}{h_t} \cdot \ell h_t)}$$

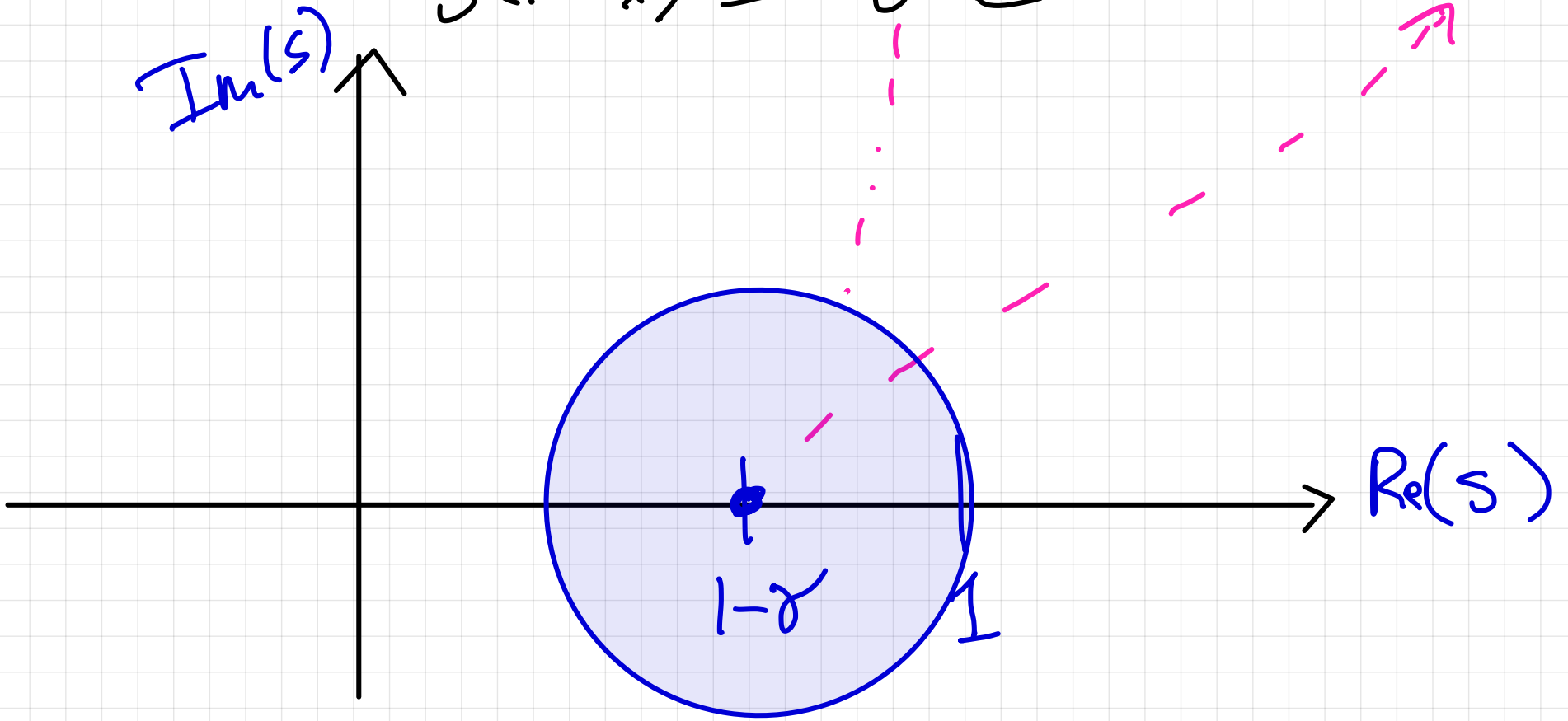
$$= z_0 e^{|\ln|s| \cdot \ell} e^{i(k_j h_x - \frac{-\phi}{h_t} \cdot \ell h_t)}$$

numerical phase velocity: $v_{ph} = \frac{-\phi(k_h x)}{k h_t}$
if not constant in k , then dispersive.

ETBS

$$\gamma = a \frac{h_t}{h_x}$$

$$S(kh_x) = \gamma e^{-ikh_x} + 1 - \gamma$$



Dispersion? ETBS

let kh_x be small

$$\Rightarrow e^{-ikh_x} \approx 1 - ikh_x$$

$$\Rightarrow s(kh_x) = 1 - \gamma + \gamma e^{-ikh_x}$$

$$\approx 1 - \gamma + \gamma(1 - ikh_x) = 1 - ikh_x$$

Also $e^{-i\omega h t} \approx 1 - i\omega h t$

$$\Rightarrow \omega h_x \approx \gamma k h_x$$

$$\Rightarrow \omega \approx \frac{\gamma k h_x}{h_x} = \frac{\gamma h_x}{h_x} \cdot \frac{k h_x}{h_x}$$

$$\approx \gamma c k$$

Conservation Laws

$$1D: \quad \partial_t u + (f(u))_x = 0$$

$$ND: \quad \partial_t u + \nabla \cdot (f(u)) = 0$$