## CS556 Iterative Methods Homework 1 Solution.

(1a) For the 1D finite difference problem on [0, L] = [0, 1] with uniform grid-spacing  $\Delta x = 1/(n+1)$ , the governing system matrix is

$$A = \frac{1}{\Delta x^2} \operatorname{tridiag}(-1, 2, -1).$$
 (1)

Since A is SPD, the largest eigenvalue satisfies

$$\lambda_{\max} \geq \max_{\underline{z} \in \mathbb{R}^n} \frac{\underline{z}^T A \underline{z}}{\underline{z}^T \underline{z}}.$$
 (2)

Take  $\underline{z} = \underline{e}_j$ , the *j*th column of *I*, the  $n \times n$  identity matrix. Then  $\underline{z}^T A \underline{z} = a_{jj} = 2/\Delta x^2 = 2(n+1)^2 \sim 2n^2$ , so,

$$\lambda_{\max}(A) \geq \frac{2}{\Delta x^2} \sim 2n^2 = O(n^2).$$
(3)

This lower bound is within a factor of two of the actual asymptotic value,  $\lambda_{\max}(A) \sim 4/\Delta x^2$ , but it is a reasonable first cut at an approximation. As we will see below, it works well for the nonuniform mesh case where a closed-form approximation is not readily accessible.

For every useful discretization of the continuous problem, we have

$$\lambda_{\min}(A) \sim \pi^2.$$
 (4)

Therefore, the condition number of A satisfies

$$\kappa(A) = \frac{\lambda_{\max}}{\lambda_{\min}} = \frac{O(n^2)}{\pi^2} = O(n^2).$$
(5)

(1b) For the same problem with nonuniform spacing, the mininum eigenvalue relationship (4) holds, so  $\lambda_{\min} \sim \pi^2$ . To get the maximum eigenvalue, look at the *j*th equation:

$$-\frac{d^2u}{dx^2}\Big|_{x_j} \approx -\frac{1}{\overline{\Delta x}} \left[\frac{u_{j+1} - u_j}{\Delta x_{j+1}} - \frac{u_j - u_{j-1}}{\Delta x_j}\right] = \frac{1}{\Delta x_j \Delta x_{j+1}} \left(\alpha u_{j-1} + \beta u_j + \gamma u_{j+1}\right), \quad (6)$$

with

$$\alpha = -\frac{\Delta x_{j+1}}{\overline{\Delta x}}, \quad \beta = 2, \quad \gamma = -\frac{\Delta x_j}{\overline{\Delta x}}.$$
(7)

Therefore, the diagonal entry is

$$a_{jj} = \frac{2}{\Delta x_j \Delta x_{j+1}} = O(\Delta x_j^{-2}), \qquad (8)$$

where the last equality holds by virtue of the assumption that the variation between  $\Delta x_j$  and  $\Delta x_{j+1}$  is bounded by a constant. (For the Chebyshev point distribution that ratio turns out to be  $\approx 3$ , rather than 2 as hinted at in the assignment.) From Part (1a), we have that

$$\lambda_{\max}(A) \geq \max_{j} |a_{jj}| = O(\Delta x_{\min}^{-2}).$$
(9)

Therefore, the condition number of A satisfies

$$\kappa(A) = \frac{\lambda_{\max}}{\lambda_{\min}} = \frac{O(\Delta x_{\min}^{-2})}{\pi^2} = O(\Delta x_{\min}^{-2}).$$
(10)

## (1c) Use Gershgorin to localize the eigenvalues of $D^{-1}A$ for both the uniform and nonuniform cases.

• Uniform Case. Here,  $D^{-1} = \frac{h^2}{2}I$ , so  $D^{-1}A = \text{tridiag}(-\frac{1}{2}, 1, -\frac{1}{2})$ . Consequently, all Gershgorin discs are centered at (1,0) in the complex  $\lambda$ -plane. On every row, save for the first and last, the disc radius is

$$\rho_i = \sum_{j \neq i} |\alpha_{ij}| = \frac{1}{2} + \frac{1}{2} = 1, \qquad (11)$$

where  $\alpha_{ij}$  are taken to be the entries of  $D^{-1}A$ . The discs associated with the first and last rows of the matrix have radius  $\frac{1}{2}$ , which implies that they lie within the union of the remaining discs and therefore do not enlarge potential domain of the eigenvalues. Thus, the eigenvalues of  $D^{-1}A$  are in the disc of radius 1 centered at (0,1) in the complex plane. Given that  $D^{-1}A$  is also real symmetric, then the eigenvalues are real and we can assert that  $\lambda \in [0, 2]$ 

• Nonuniform Case. From (6)–(7) we see that  $D^{-1}A$  for the nonuniform case leads to

$$\left[D^{-1}A\right]_{j} = \operatorname{tridiag}\left(\frac{\alpha}{\beta}, 1, \frac{\gamma}{\beta}\right). \tag{12}$$

By construction, the Gershgorin discs are all centered at (1,0). For the interior points, the radii are

$$\rho_j = \left|\frac{\alpha}{\beta}\right| + \left|\frac{\gamma}{\beta}\right| + = \frac{1}{2\overline{\Delta x}}\left(|\Delta x_{j+1}| + |\Delta x_j|\right) = 1.$$
(13)

As in the case of the uniform points, the discs associated with the end points are centered at (1,0) and within the unit discs of the remaining rows.

The preceding analysis shows that diagonal scaling successfully localizes the eigenvalues for both the uniform and nonuniform grid spacing cases.

Unfortunately, this is a necessary condition but in this case not sufficient to prove convergence because the largest Gershgorin disc touches both endpoints of the interval [0,2], which means that the spectral radius of  $E = I - D^{-1}A$  could be as large as 1, rather than strictly < 1. This case is unlike the Helmholtz problem considered in class,  $H = I + \gamma^2 A$ , where Gershgorin would immediately localize the discs of E to have radii < 1, independent of n.



Figure 1: Chebyshev points on [-1,1].

- (1d) Figure 1 shows the Chebyshev points,  $z_j = cos(\theta_j)$  for uniformly-distributed angles  $\theta_0, \ldots, \theta_{n+1}$  on  $[-\pi, \pi]$ . (Our points are translated and scaled to [0, 1],  $x_j = \frac{1}{2}(1 + z_j)$ .) Here, we discuss the relative spacings, per the HW1 questions.
  - As illustrated in Fig. 1, the maximum spacing  $\Delta z_j \sim \frac{\pi}{n+1}$ , which corresponds to the arclength of each segment. As  $\Delta \theta_j \longrightarrow 0$  the distance between the center points tends towards this arclength. Similarly, max  $\Delta x_j \sim \frac{\pi}{2(n+1)}$ .
  - From Fig. 1 scales as

$$\Delta z_{\min} \sim 1 - \cos \frac{\pi}{n+1} \sim \frac{\theta^2}{2}, \qquad (14)$$

where  $\theta := \pi/(n+1)$ . Consequently,  $\Delta x_{\min} \sim \theta^2/4$ .

• The maximum ratio occurs near the ends of the domain. After the first point, the second grid point is located at  $x_2 \sim (2\theta)^2/4$ , so that the second space is of size  $\sim 3\theta^2/4 \sim 3\Delta x_{\min}$ . Thus, the slowly varying condition holds (adjacent spacings vary by at most a factor of 3).

In fact, a closer inspection reveals that

$$a_{jj} = \frac{2}{\Delta x_j \Delta x_{j+1}}.$$
(15)

Since the size ratio for the smallest Chebyshev points is 3, we can anticipate that  $\max_j a_{jj} \sim 2/(3\Delta x_{min}^2)$ , which is plotted in Fig. 2, along with the actual eigenvalues.

**Summary.** The maximum eigenvalue and condition number of A for the uniform case scale like  $O(n^2)$  while those for the Chebyshev case scale like  $O(n^4)$ , which is much more severe. Diagonal scaling, however, brings all the eigenvalues into the unit disc for both cases and yields a condition number that scales as  $O(n^2)$ , as illustrated in Fig. 1.



Figure 2: Computed (symbols) and modeled (lines) eigenvalues for Poisson operator on Chebyshev points: green= $\lambda_{\text{max}}$ , blue= $\lambda_{\text{min}} \sim \pi^2$ . Also shown is the condition number for  $A_{\text{cheb}}$ , which scales as  $O(n^4)$ , and for  $D^{-1}A_{\text{cheb}}$ , which scales as  $O(n^2)$ .

- 2. With each power of the tridiagonal matrix A the number of diagonals increases by 2. So we start with ~ 3n nonzeros, then ~ 5n for k = 2, and so forth. Neglecting end effects (i.e., for  $k \ll n$ ) the number of nonzeros for  $A^k$  is ~ n(1+2k).
- **3.** Consider a vector  $\underline{e}$  that is zero everywhere except for entries  $e_j$ ,  $j = j_0 : j_1$ . The matrix-vector product  $\underline{w} = A\underline{e}$  will be a linear combination of columns  $\underline{a}_{j_0} : \underline{a}_{j_1}$  of A, which will have nonzeros in locations  $j = j'_0 : j'_1$ , where  $j'_0 = \max(1, j_0 1)$  and  $j'_1 = \min(n, j_1 + 1)$ . Thus, if we start with  $\underline{e} = \underline{e}_j$ , the *j*th column of the identity matrix, the number of nonzeros will grows by 1 on either side of *j* for each matrix-vector product. If j = n/2 (assuming *n* even), then it will take k = n/2 matrix-vector products to produce a nonzero in every location of  $A^k \underline{e}_j$ .
- 4. Figure 3 shows the columns of  $Z := A^{-1}$ . We see that all the entries are strictly positive and that  $A^{-1}$  is consequently completely full.

If we were to solve  $A\underline{u} = \underline{f}$  problem in parallel, with  $\underline{f}$  a distributed data vector and  $\underline{u}$  a distributed solution vector then we would need an all-to-all communication because *each* nonzero in  $\underline{f}$  has a nontrivial influence on every value of  $\underline{u}$ . This assertion is a direct consequence of the fact that every column vector of  $A^{-1}$  is completely full.



Figure 3: Columns of  $A^{-1}$  for n=20, where A is the 1D Poisson operator on a uniform grid.