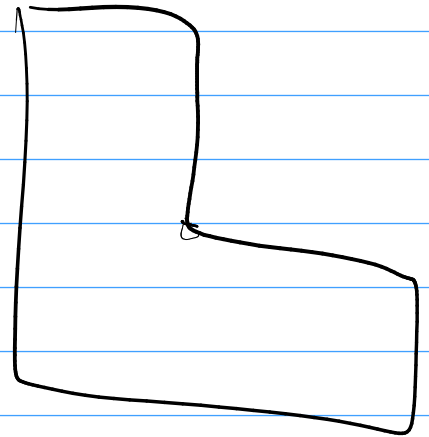


TODAY:

- Obdwi: \wedge
- Helmholtz: she bad parks 5
- Calderón
- numerice

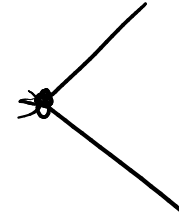
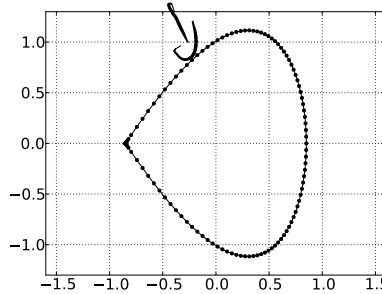
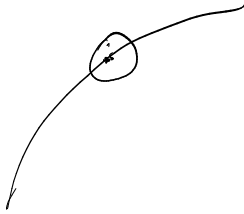


$$\langle u, v \rangle = \int u(x)v(x) dx \quad \int_{\mathcal{R}} u \Delta v - v \Delta u = O(h) \sim S(h \cdot \mathcal{D}1)$$

$v = G$ \mathcal{C}^2

Domains with Corners

$[-1, 1]$



What's the problem? (*Hint: Jump condition for constant density*)

At corner x_0 : (2D)

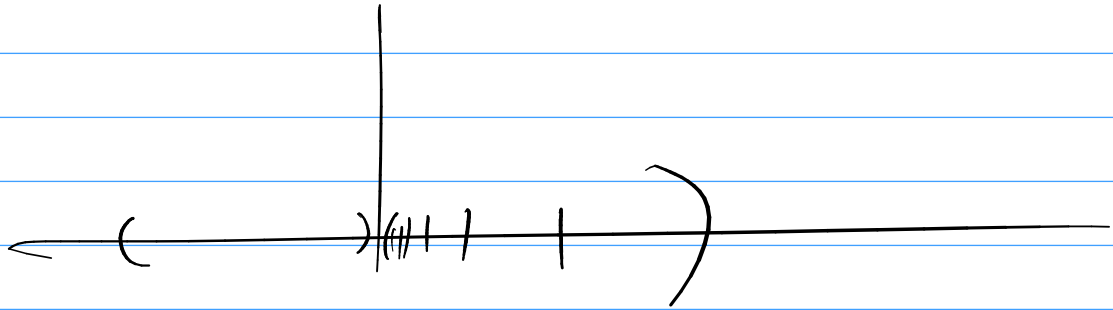
$$\lim_{x \rightarrow x_0 \pm} = \int_{\partial\Omega} \hat{n} \cdot \nabla_y G(x, y) \phi(y) ds_y \pm \frac{1}{2} \frac{\langle \text{opening angle on } \pm \text{ side} \rangle}{\pi} \phi$$

→ non-continuous behavior of potential on Γ at x_0

What space have we been living in?

Fixes:

- $I +$ Bounded (Neumann) + Compact (Fredholm)



1
x

- Use L^2 theory

(point behavior “invisible”)

Numerically: Needs consideration, but ultimately easy to fix.

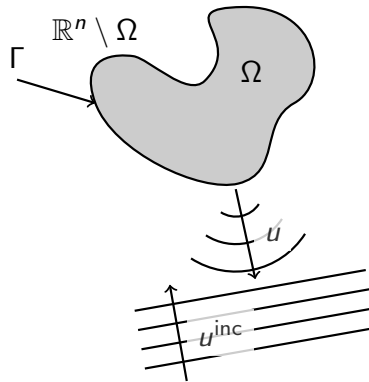
8.2 Helmholtz

Where does Helmholtz come from?

Derive the Helmholtz equation from the wave equation

$$\partial_t^2 U = c^2 \Delta U,$$

The prototypical Helmholtz BVP: A Scattering Problem



Ansatz:

$$u^{\text{tot}} = u + u^{\text{inc}}$$

Solve for scattered field u .

Helmholtz: Some Physics

Physical quantities:

- Velocity potential: $U(x, t) = u(x)e^{-i\omega t}$
(fix phase by e.g. taking real part)
- Velocity: $v = (1/\rho_0)\nabla U$
- Pressure: $p = -\partial_t U = i\omega u e^{-i\omega t}$
 - Equation of state: $p = f(\rho)$

What's ρ_0 ?

What happens to a pressure BC as $\omega \rightarrow 0$?

Helmholtz: Boundary Conditions

- **Sound-soft:** Pressure remains constant
 - Scatterer “gives”
 - $u = f \rightarrow$ Dirichlet
- **Sound-hard:** Pressure same on both sides of interface
 - Scatterer “does not give”
 - $\hat{n} \cdot \nabla u = 0 \rightarrow$ Neumann
- **Impedance:** Some pressure translates into motion
 - Scatterer “resists”
 - $\hat{n} \cdot \nabla u + ik\lambda u = 0 \rightarrow$ Robin ($\lambda > 0$)
- **Sommerfeld** radiation condition: allow only outgoing waves

$$r^{\frac{n-1}{2}} \left(\frac{\partial}{\partial r} - ik \right) u(x) \rightarrow 0 \quad (r \rightarrow \infty)$$

(where n is the number of space dimensions)

Many interesting BCs → many IEs! :)

Transmission between media: What's continuous?

Unchanged from Laplace

Theorem 18 (Green's Formula) [Colton/Kress IAEST Thm 2.1] If $\Delta u + k^2 u = 0$, then

$$(S(\hat{n} \cdot \nabla u) - Du)(x) = \begin{cases} u(x) & x \in D \\ \frac{u(x)}{2} & x \in \partial D \\ 0 & x \notin D \end{cases}$$

$$\begin{aligned} \lim_{x \rightarrow x_0 \pm} (S'u) &= \left(S' \mp \frac{1}{2} I \right) (u)(x_0) \Rightarrow [Su] = 0 \\ \lim_{x \rightarrow x_0 \pm} (Du) &= \left(D \pm \frac{1}{2} I \right) (u)(x_0) \Rightarrow [S'u] = -u \\ &= 0 \end{aligned}$$

$M_0^{(1)}(k \cdot r)$

Why is singular behavior (esp. jump conditions) unchanged?

$$\frac{e^{ikr}}{r} = \frac{1}{r} + \frac{e^{ikr} - 1}{r}$$

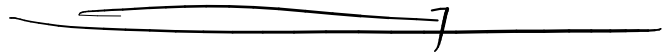
Why does Green's formula survive?

Remember Green's theorem:

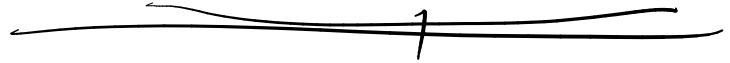
$$\int_{\Omega} u \Delta v - v \Delta u = \int_{\partial \Omega} u(\hat{n} \cdot \nabla v) - v(\hat{n} \cdot \nabla u) ds$$

$$\Delta u + k^2 u = 0$$

' Δ



$\Delta + k^2$



Resonances

— Δ on a bounded (interior) domain with homogeneous Dirichlet/Neumann BCs has countably many real, positive eigenvalues.

What does that have to do with Helmholtz?

Why could it cause grief?

Helmholtz: Boundary Value Problems

Find $u \in C(\bar{D})$ with $\Delta u + k^2 = 0$ such that

	Dirichlet	Neumann
Int.	$\lim_{x \rightarrow \partial D^-} u(x) = g$ ⊕ unique (-resonances)	$\lim_{x \rightarrow \partial D^-} \hat{n} \cdot \nabla u(x) = g$ ⊕ unique (-resonances)
Ext.	$\lim_{x \rightarrow \partial D^+} u(x) = g$ Sommerfeld ⊕ unique	$\lim_{x \rightarrow \partial D^+} \hat{n} \cdot \nabla u(x) = g$ Sommerfeld ⊕ unique

with $g \in C(\partial D)$.

Find layer potential representations for each.

$$u(x) = S\sigma - iD(\sigma)$$

"CFIE"

*inherits nullspaces
(because adj.)
spurious resonances*

Combined-field ext Neumann; $\left(\frac{I}{2} - S' - D'\right)\sigma = g$

$$c(\lambda) = \int \sigma - i D \int \sigma$$

$$D'S = (S')^2 - \frac{I}{4}$$

Patching up resonances

Issue: Ext. IE inherits non-uniqueness from 'adjoint' int. BVP

Fix: Tweak representation [Brakhage/Werner '65, ...]

(also called the 'CFIE' – 'combined field integral equation')

$$u = D\phi - i\alpha S\phi$$

(α : tuning knob $\rightarrow 1$ is fine, $\sim k$ better for large k)

How does this help?

Uniqueness for remaining IEs similar. (skipped)

8.3 Calderón identities

Show that D' is self adjoint.

Show that $(S\varphi, D'\psi) = ((S' + I/2)\varphi, (D - I/2)\psi)$.


$(\varphi, SD'\psi)$?

$$(D'\varphi, \psi) = (\varphi, D'\psi) \quad |$$

$$u := D\varphi \quad v := D\psi$$

$$\int_{\Omega} u \Delta v - v \Delta u = \int_{\partial\Omega} (n \cdot \nabla u) v - u (n \cdot \nabla v) = 0$$

$$\begin{aligned}
(D' \varphi, \psi) &= (D' \varphi, [v]) \\
&= (D' \varphi, v^+) - (D' \varphi, v^-) \\
&= ((D\varphi)^+, v') - ((D\varphi)^-, v') \\
&= ([u], v') = (\varphi, D' \psi)
\end{aligned}$$



$$S D' \psi = \left(D^2 - \frac{\mp}{\varphi} \right) \psi$$

Calderón Identities: Summary

- $SD' = D^2 - I/4$
- $D'S = S'^2 - I/4$

Also valid for Laplace (jump relation same after all!)

Why do we care?

9 Back from Infinity: Discretization

9.1 Fundamentals: Meshes, Functions, and Approximation

Numerics: What do we need?

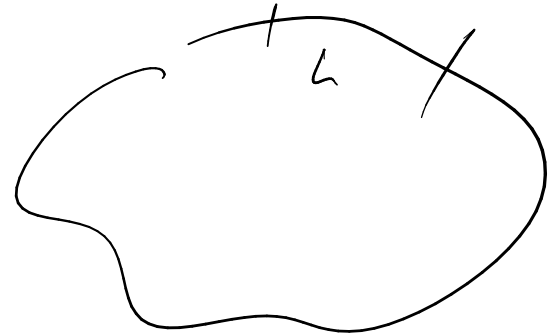
- Discretize curves and surfaces
 - Interpolation
 - Grid management
 - Adaptivity
- Discretize densities
- Discretize integral equations
 - Nyström, Collocation, Galerkin
- Compute integrals on them
 - “Smooth” quadrature
 - Singular quadrature
- Solve linear systems

$$A u = f$$
$$(A u, \varphi) = (f, \varphi)$$

Constructing Discrete Function Spaces

Floating point numbers → Functions
(Degrees of Freedom/DoFs) ←

$$\|f - f_h\| \leq C \cdot h$$



Discretization relies on three things:

- Base/reference domain
- Basis of functions
- Meaning of DoFs

Related finite element concept: *Ciarlet triple*

Discretization options for a curve?

What do the DoFs mean?

Common DoF choices:

- Point values of function
- Point values of (directional?) derivatives
- Basis coefficients
- Moments

Often: useful to have both “modes”, “nodes”, jump back and forth

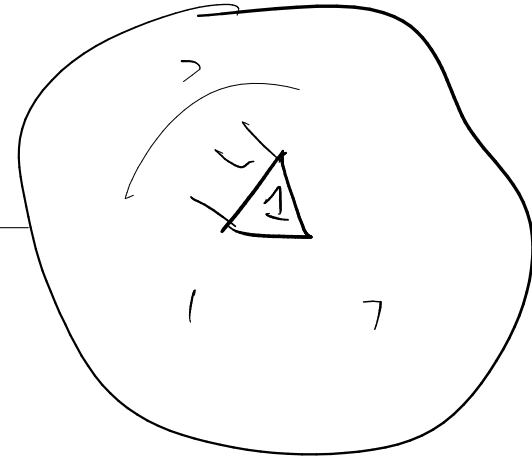
Why high order?

Order p : Error bounded as

$$|u_h - u| \leq Ch^p$$

Thought experiment:

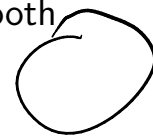
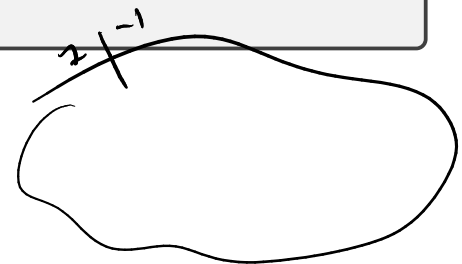
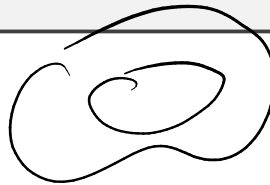
First order	Fifth order
1,000 DoFs \approx 1,000 triangles Error: 0.1	1,000 DoFs \approx 66 triangles Error: 0.1
Error: 0.01 \rightarrow ? <i>100,000 triangles</i>	Error: 0.01 \rightarrow ?



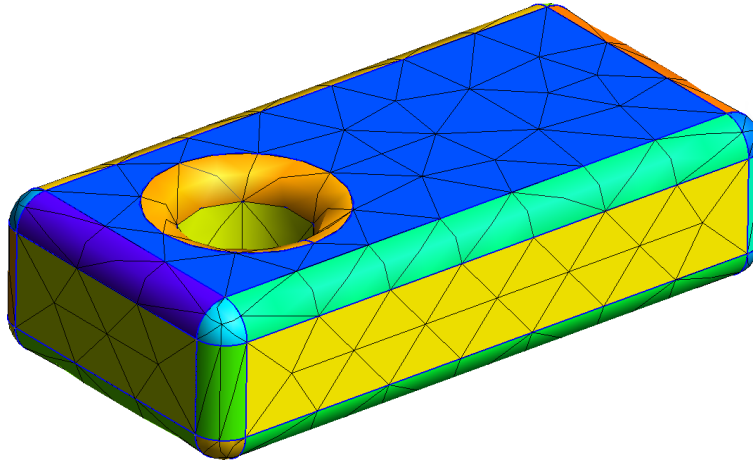
Complete the table.

Remarks:

- Want $p \geq 3$ available.
- **Assumption:** Solution sufficiently smooth.
- Ideally: p chosen by user



What is an Unstructured Mesh?



Why have an unstructured mesh?

What is the trade-off in going unstructured?

Fixed-order vs Spectral

Fixed-order	Spectral
Number of DoFs n	Number of DoFs n
\sim	\sim
Number of 'elements'	Number of modes resolved
Error $\sim \frac{1}{n^p}$	Error $\sim \frac{1}{C^n}$
Examples?	Examples?
<ul style="list-style-type: none">• Piecewise Polynomials	<ul style="list-style-type: none">• Global Fourier• Global Orth. Polynomials

What assumptions are buried in each of these?

What should the DoFs be?

What's the difficulty with purely modal discretizations?

Vandermonde Matrices

$$\begin{pmatrix} x_0^0 & x_0^1 & \cdots & x_0^n \\ x_1^0 & x_1^1 & \cdots & x_1^n \\ \vdots & \vdots & \ddots & \vdots \\ x_n^0 & x_n^1 & \cdots & x_n^n \end{pmatrix} \begin{pmatrix} a_0 \\ a_1 \\ \vdots \\ a_n \end{pmatrix} = ?$$

Generalized Vandermonde Matrices

$$\begin{pmatrix} \phi_0(x_0) & \phi_1(x_0) & \cdots & \phi_n(x_0) \\ \phi_0(x_1) & \phi_1(x_1) & \cdots & \phi_n(x_1) \\ \vdots & \vdots & \ddots & \vdots \\ \phi_0(x_n) & \phi_1(x_n) & \cdots & \phi_n(x_n) \end{pmatrix} \begin{pmatrix} a_0 \\ a_1 \\ \vdots \\ a_n \end{pmatrix} = ?$$

Generalized Vandermonde Matrices

$$\begin{pmatrix} \phi_0(x_0) & \phi_1(x_0) & \cdots & \phi_n(x_0) \\ \phi_0(x_1) & \phi_1(x_1) & \cdots & \phi_n(x_1) \\ \vdots & \vdots & \ddots & \vdots \\ \phi_0(x_n) & \phi_1(x_n) & \cdots & \phi_n(x_n) \end{pmatrix} \text{MODAL COEFFS} = \text{NODAL COEFFS}$$

Node placement?

Vandermonde conditioning?

What about multiple dimensions?

Common Operations

(Generalized) Vandermonde matrices simplify common operations:

- Modal \leftrightarrow Nodal (“Global interpolation”)
 - Filtering
 - Up-/Oversampling
- Point interpolation (Hint: solve using V^T)
- Differentiation
- Indefinite Integration
- Inner product
- Definite integration