

Today

- Exh Dirichlet
- Domains w/ Γ
- Helmholtz

Announcements

- Projects / presentations

Uniqueness of Integral Equation Solutions

Theorem (Nullspaces [Kress LIE 2nd ed. Thm 6.20])

- ▶ $N(I/2 - D) = N(I/2 - S') = \{0\}$
- ▶ $N(I/2 + D) = \text{span}\{1\}$, $N(I/2 + S') = \text{span}\{\psi\}$,
where $\int \psi \neq 0$.

IE Uniqueness: Proofs (1/3)

Show $N(I/2 - D) = \{0\}$.



IE Uniqueness: Proofs (2/3)

Show $N(I/2 - S') = \{0\}$.



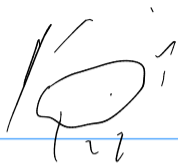
IE Uniqueness: Proofs (3/3)

Show $N(I/2 + D) = \text{span}\{1\}$.

What conditions on the RHS do we get for int. Neumann and ext. Dirichlet?

→ “Clean” Existence for 3 out of 4.

Ext. Dirichlet



$$u(x) = \tilde{D}\sigma(x)$$

$$G(x,y) = \hat{n} \cdot \nabla_y G(x,y) + 1$$

$$\frac{1}{|x|^{n-2}} \int_{\partial\Omega} \sigma$$

constant

$$\int_{\partial\Omega} \frac{1}{|x|^{n-2}} \sigma(y) dy$$

$$\tilde{D}\sigma(\vec{x}) = \int_{\partial\Omega} \hat{n} \cdot \nabla_y G(x,y) \sigma(y) dy + \int_{\partial\Omega} \frac{1}{|x|^{n-2}} \sigma(y) dy$$

Patching up Exterior Dirichlet

Problem: $N(I/2 + S') = \{\psi\}$ do not know ψ . Use different kernel:

$$\hat{n} \cdot \nabla_y G(x, y) \quad \rightarrow \quad \hat{n} \cdot \nabla_y G(x, y) + \frac{1}{|x|^{n-2}}$$

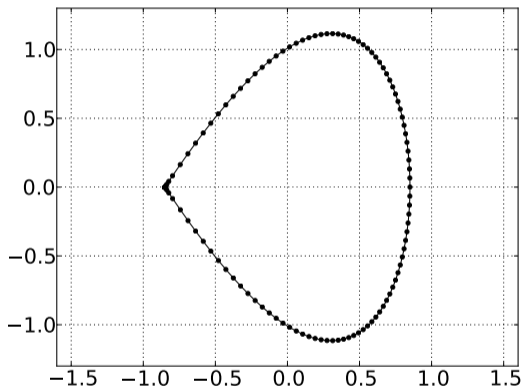
Note: Singularity only at origin! (assumed $\in \Omega$)

- ▶ 2D behavior? 3D behavior?
- ▶ Still a solution of the PDE? Compact?
- ▶ Jump condition? Exterior limit? Deduce $u = 0$ on exterior.
 - ▶ Consider $\partial_n G = O(1/r^{n-1})$.
- ▶ $|x|^{n-2} u(x) = ?$ as $|x| \rightarrow \infty$?
- ▶ Thus $\int \phi = 0$. Contribution of the second term?
- ▶ $\phi/2 + D\phi = 0$, i.e. $\phi \in N(I/2 + D) = \text{span}\{1\} \Rightarrow \psi = 0$ -
- ▶ Existence/uniqueness?

→ Existence for 4 out of 4.

Domains with Corners

$$01 = -1_{\Omega}$$



What's the problem? (*Hint: Jump condition for constant density*)

Domains with Corners (II)



At corner x_0 : (2D)

$$\lim_{x \rightarrow x_0 \pm} = \int_{\partial\Omega} \hat{n} \cdot \nabla_y G(x, y) \phi(y) ds_y \pm \frac{1}{2} \frac{\langle \text{opening angle on } \pm \text{ side} \rangle}{\pi} \phi$$

→ non-continuous behavior of potential on Γ at x_0

What space have we been living in?



Fixes:

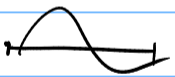
- ▶ $I +$ Bounded (Neumann) + Compact (Fredholm)
- ▶ Use L^2 theory
(point behavior "invisible")

Numerically: Needs consideration, can drive up cost through refinement.

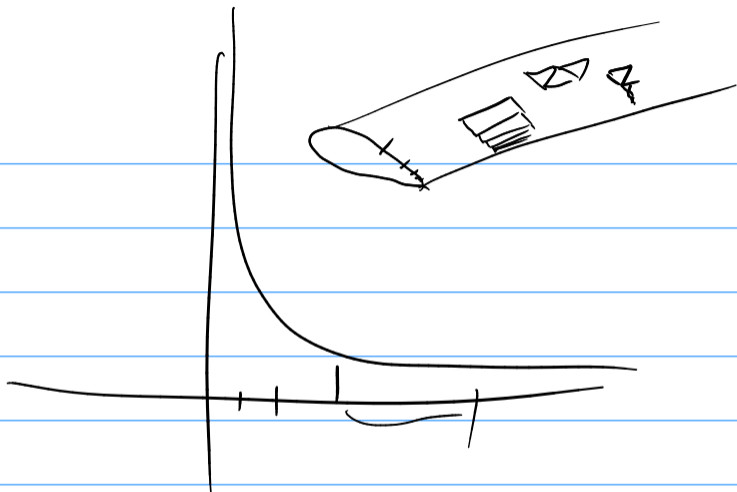


$$L^2(\Omega) = \left\{ \phi : \int_{\Omega} |\phi|^2 < \infty \right\} \quad H^1(\Omega)$$

H' :  ✓
 H' :  X ✓



$$S: H^{-1/2} \rightarrow H^{1/2}$$



Where does Helmholtz come from?

$$u(t, x, y)$$

Derive the Helmholtz equation from the wave equation $\partial_t^2 U = c^2 \Delta U$. Q:
What is c ?

$$u(x, t) = e^{-i\omega t} u(x)$$

$$u(x) \partial_t^2 (e^{-i\omega t}) = c^2 e^{-i\omega t} \Delta u$$

$$u(x) (-i\omega)^2 e^{-i\omega t} = c^2 e^{-i\omega t} \Delta u$$

$$-\omega^2 u(x) = c^2 \Delta u$$

$$u(x) = \frac{c^2}{\omega^2} \Delta u$$
$$k^2$$

Helmholtz vs. Yukawa

Helmholtz Equation

- ▶ $\Delta u + k^2 u(x) = 0$
- ▶ Indefinite operator
- ▶ Oscillatory solution
- ▶ Difficult to solve, especially for large k

$$\frac{e^{-ikr}}{r}$$



Yukawa Equation

- ▶ $-\Delta u \oplus k^2 u(x) = 0$
- ▶ Positive definite operator
- ▶ Smooth solutions
- ▶ 'Screened Coulomb' interaction
- ▶ Generally quite simple to solve



Helmholtz vs. Yukawa

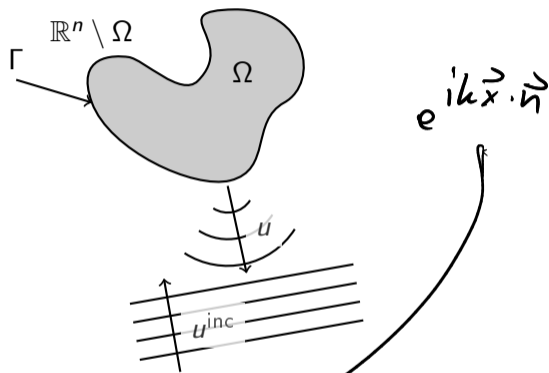
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The prototypical Helmholtz BVP: A Scattering Problem



Ansatz:

$$u^{tot} = (u) + u^{inc}$$

Solve for scattered field u .

Helmholtz: Some Physics

Colton/Kress

Physical quantities:

- ▶ Velocity potential: $U(x, t) = u(x)e^{-i\omega t}$
(fix phase by e.g. taking real part)
- ▶ Velocity: $v \in (1/\rho_0)\nabla U$
- ▶ Pressure: $p = -\partial_t U = i\omega u e^{-i\omega t}$
 - ▶ Equation of state: $p = f(\rho)$

"sound-hard"
"sound-soft"

What's ρ_0 ?

density about which we've linearized

What happens to a pressure BC as $\omega \rightarrow 0$?

Helmholtz: Boundary Conditions

Interfaces between media: What's continuous?

- ▶ **Sound-soft:** Scatterer “gives”
 - ▶ Pressure remains constant in time
 - ▶ $u = f \rightarrow$ Dirichlet
- ▶ **Sound-hard:** Scatterer “does not give”
 - ▶ Pressure varies, same on both sides of interface
 - ▶ $\hat{n} \cdot \nabla u = 0 \rightarrow$ Neumann
- ▶ **Impedance:** Some pressure translates into motion
 - ▶ Scatterer “resists”
 - ▶ $\hat{n} \cdot \nabla u + ik\lambda u = 0 \rightarrow$ Robin ($\lambda > 0$)
- ▶ **Sommerfeld** radiation condition: allow only outgoing waves (n -dim)

$$r^{\frac{n-1}{2}} \left(\frac{\partial}{\partial r} - ik \right) u(x) \rightarrow 0 \quad (r \rightarrow \infty)$$

Many interesting BCs \rightarrow many IEs! :)

$$0 = (\partial_t^2 - \partial_x^2) u = \underbrace{(\partial_t - \partial_x)^{-i\omega}}_{-i\omega} (\partial_t + \partial_x)^{-i\omega} u$$

$$(\partial_t - \partial_x) u = 0 \quad \text{"advection"}$$

$$u(x, t) = \bar{u}(t + x)$$

Unchanged from Laplace

Theorem (Green's Formula [Colton/Kress IAEST Thm 2.1])

If $\Delta u + k^2 u = 0$, then

$$(S(\hat{n} \cdot \nabla u) - Du)(x) = \begin{cases} u(x) & x \in D \\ \frac{u(x)}{2} & x \in \partial D \\ 0 & x \notin D \end{cases}$$

$$\begin{aligned} \lim_{x \rightarrow x_0 \pm} (S'u) &= \left(S' \mp \frac{1}{2} I \right) (u)(x_0) &\Rightarrow [Su] = 0 \\ \lim_{x \rightarrow x_0 \pm} (Du) &= \left(D \pm \frac{1}{2} I \right) (u)(x_0) &\Rightarrow [S'u] = -u \\ & &\Rightarrow [Du] = u \\ & &\Rightarrow [D'u] = 0 \end{aligned}$$

Unchanged from Laplace

Why is singular behavior (esp. jump conditions) unchanged?

$$\frac{e^{ikr}}{r} = \frac{1}{r} + o\left(\frac{1}{r}\right) \quad \text{as } r \rightarrow 0$$

Why does Green's formula survive?

$$\int_{\Omega} u \Delta v - v \Delta u = \int_{\partial \Omega} u(\hat{n} \cdot \nabla v) - v(\hat{n} \cdot \nabla u)$$

$$u \rightarrow \Delta u + k^2 u = 0$$

$$v \rightarrow G$$

Resonances

$$\lambda \neq \lambda_n > 0$$

$$\int_{\Omega} \Delta u v = \int_{\Omega} u \Delta v$$
$$\int_{\Omega} u \Delta u \geq 0$$

— Δ on a bounded (interior) domain with homogeneous Dirichlet/Neumann BCs has countably many real, positive eigenvalues.

What does that have to do with Helmholtz?

$$-\Delta u = \lambda u \quad (u=0)$$
$$\Delta u + k^2 u = 0 \quad (\lambda > 0)$$

Why could it cause grief?

inf. BVPs ; non-uniqueness

Helmholtz: Boundary Value Problems

Find $u \in C(\bar{D})$ with $\Delta u + k^2 = 0$ such that

	Dirichlet	Neumann
Int.	$\lim_{x \rightarrow \partial D^-} u(x) = g$ ⓪ unique (−resonances)	$\lim_{x \rightarrow \partial D^-} \hat{n} \cdot \nabla u(x) = g$ ⓪ unique (−resonances)
Ext.	$\lim_{x \rightarrow \partial D^+} u(x) = g$ Sommerfeld ⊕ unique	$\lim_{x \rightarrow \partial D^+} \hat{n} \cdot \nabla u(x) = g$ Sommerfeld ⊕ unique

with $g \in C(\partial D)$.

Find layer potential representations for each.

int. Dirichlet: $\frac{\mathbb{I}}{2} - D$

ext. Dirichlet: $\frac{\mathbb{I}}{2} + D$

int. Neumann: $\frac{\mathbb{I}}{2} + S'$

Patching up resonances

Issue: Ext. IE inherits non-uniqueness from 'adjoint' int. BVP

Fix: Tweak representation [Brakhage/Werner '65, ...]
(also called the *CFIE* or *combined field integral equation*)

$$u = D\phi - \underbrace{i\alpha S\phi}$$

(α : tuning knob $\rightarrow 1$ is fine, $\sim k$ better for large k)